

Having defined the unrestricted density matrices \mathbf{P}^α , \mathbf{P}^β , \mathbf{P}^T , and \mathbf{P}^S we will now use these definitions to give explicit form to the unrestricted Fock matrices F^α and F^β .

3.8.4 Expression for the Fock Matrices

To obtain expressions for the elements of the matrices F^α and F^β , we simply take matrix elements in the basis $\{\phi_\mu\}$ of the two Fock operators f^α (Eq. (3.316)) and f^β (Eq. (3.318)), and use expressions (3.322) to (3.326) for matrix elements of the coulomb and exchange operators. That is,

$$\begin{aligned} F_{\mu\nu}^\alpha &= \int d\mathbf{r}_1 \phi_\mu^*(1) f^\alpha(1) \phi_\nu(1) \\ &= H_{\mu\nu}^{\text{core}} + \sum_a^{N^\alpha} [(\phi_\mu \phi_\nu | \psi_a^\alpha \psi_a^\alpha) - (\phi_\mu \psi_a^\alpha | \psi_a^\alpha \phi_\nu)] + \sum_a^{N^\beta} (\phi_\mu \phi_\nu | \psi_a^\beta \psi_a^\beta) \end{aligned} \quad (3.346)$$

$$\begin{aligned} F_{\mu\nu}^\beta &= \int d\mathbf{r}_1 \phi_\mu^*(1) f^\beta(1) \phi_\nu(1) \\ &= H_{\mu\nu}^{\text{core}} + \sum_a^{N^\beta} [(\phi_\mu \phi_\nu | \psi_a^\beta \psi_a^\beta) - (\phi_\mu \psi_a^\beta | \psi_a^\beta \phi_\nu)] + \sum_a^{N^\alpha} (\phi_\mu \phi_\nu | \psi_a^\alpha \psi_a^\alpha) \end{aligned} \quad (3.347)$$

To continue, we substitute the basis set expansions of ψ_a^α and ψ_a^β to get

$$\begin{aligned} F_{\mu\nu}^\alpha &= H_{\mu\nu}^{\text{core}} + \sum_\lambda \sum_\sigma \sum_a^{N^\alpha} C_{\lambda a}^\alpha (C_{\sigma a}^\alpha)^* [(\mu\nu | \sigma\lambda) - (\mu\lambda | \sigma\nu)] + \sum_\lambda \sum_\sigma \sum_a^{N^\beta} C_{\lambda a}^\beta (C_{\sigma a}^\beta)^* (\mu\nu | \sigma\lambda) \\ &= H_{\mu\nu}^{\text{core}} + \sum_\lambda \sum_\sigma P_{\lambda\sigma}^\alpha [(\mu\nu | \sigma\lambda) - (\mu\lambda | \sigma\nu)] + \sum_\lambda \sum_\sigma P_{\lambda\sigma}^\beta (\mu\nu | \sigma\lambda) \\ &= H_{\mu\nu}^{\text{core}} + \sum_\lambda \sum_\sigma P_{\lambda\sigma}^T (\mu\nu | \sigma\lambda) - P_{\lambda\sigma}^\alpha (\mu\lambda | \sigma\nu) \end{aligned} \quad (3.348)$$

$$\begin{aligned} F_{\mu\nu}^\beta &= H_{\mu\nu}^{\text{core}} + \sum_\lambda \sum_\sigma \sum_a^{N^\beta} C_{\lambda a}^\beta (C_{\sigma a}^\beta)^* [(\mu\nu | \sigma\lambda) - (\mu\lambda | \sigma\nu)] + \sum_\lambda \sum_\sigma \sum_a^{N^\alpha} C_{\lambda a}^\alpha (C_{\sigma a}^\alpha)^* (\mu\nu | \sigma\lambda) \\ &= H_{\mu\nu}^{\text{core}} + \sum_\lambda \sum_\sigma P_{\lambda\sigma}^\beta [(\mu\nu | \sigma\lambda) - (\mu\lambda | \sigma\nu)] + \sum_\lambda \sum_\sigma P_{\lambda\sigma}^\alpha (\mu\nu | \sigma\lambda) \\ &= H_{\mu\nu}^{\text{core}} + \sum_\lambda \sum_\sigma P_{\lambda\sigma}^T (\mu\nu | \sigma\lambda) - P_{\lambda\sigma}^\beta (\mu\lambda | \sigma\nu) \end{aligned} \quad (3.349)$$

If one compares these expressions with the corresponding restricted closed-shell expression (3.154), one sees that the coulomb term is identical and depends on the total density matrix. The difference is only that here one has separate representations of the α and β density matrices rather than, as in the closed-shell case,

$$P_{\mu\nu}^\alpha = P_{\mu\nu}^\beta = \frac{1}{2} P_{\mu\nu}^T \quad (3.350)$$

The coupling of the two sets of equations is made explicit in the above expressions, i.e., F^α depends on \mathbf{P}^β (through the total density matrix \mathbf{P}^T) and F^β similarly depends on \mathbf{P}^α .